

High frequency
analysis of the
dissipative Helmholtz
equation

Julien ROYER

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The Helmholtz equation

We study on \mathbb{R}^n the following Helmholtz equation:

$$(-h^2 \Delta + V_1(x) - ihV_2(x) - E)u = S.$$

This equation models accurately the propagation of the electromagnetic field of a laser in material medium.

$$\begin{aligned} V_1(x) - E & : \text{refraction index,} \\ V_2(x) & : \text{absorption index,} \\ S & : \text{source term,} \\ h & : \text{wave length,} \end{aligned}$$

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The Helmholtz equation

We study on \mathbb{R}^n the following Helmholtz equation:

$$(-h^2 \Delta + V_1(x) - ihV_2(x) - E)u_h = S.$$

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We consider the high frequency approximation $h \rightarrow 0$.

The non-selfadjoint Schrödinger operator

- When V_2 is constant, it can be put in the spectral parameter:

$$(H_1^h - z_h)u_h = S,$$

with

$$H_1^h = -h^2\Delta + V_1(x) \quad \text{and} \quad z_h = E + ihV_2.$$

The non-selfadjoint Schrödinger operator

- When V_2 is **constant**, it can be put in the **spectral parameter**:

$$(H_1^h - z_h)u_h = S,$$

with

$$H_1^h = -h^2\Delta + V_1(x) \quad \text{and} \quad z_h = E + ihV_2.$$

- When V_2 is **variable**, it has to be **in the operator itself**:

$$(H_h - E)u_h = S,$$

with

$$H_h = -h^2\Delta + V_1(x) - ihV_2(x).$$

↪ we have to work with a **non-selfadjoint** operator.

Dissipative operators

The operator H on the Hilbert space \mathcal{H} is said to be **dissipative** if

$$\forall \varphi \in \mathcal{D}(H), \quad \operatorname{Im} \langle H\varphi, \varphi \rangle \leq 0.$$

H is said to be **maximal dissipative** if any dissipative extension of H is trivial.

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- The **resolvent** $(H - z)^{-1}$ is well-defined if $\operatorname{Im} z > 0$ and

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- H generates a **contractions semi-group**

$$t \in \mathbb{R}_+ \mapsto e^{-itH}, \quad \|e^{-itH}\|_{\mathcal{L}(\mathcal{H})} \leq 1,$$

and for $\varphi \in \mathcal{D}(H)$:

$$\frac{d}{dt} \|e^{-itH} \varphi\|_{\mathcal{H}}^2 = 2 \operatorname{Im} \langle He^{-itH} \varphi, e^{-itH} \varphi \rangle \leq 0.$$

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- ① We first look for **uniform resolvent estimates**:

$$\sup_{\substack{\operatorname{Re} z \sim E \\ \operatorname{Im} z > 0}} \|(H - z)^{-1}\|_{\mathcal{L}(\mathcal{H}_1, \mathcal{H}_1^*)} \leq c$$

$$(\mathcal{H}_1 \subset L^2(\mathbb{R}^n) \subset \mathcal{H}_1^*).$$

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$(\mathcal{H}_1 \subset L^2(\mathbb{R}^n) \subset \mathcal{H}_1^*)$. This gives the **limiting absorption principle**:

$$\lim_{\mu \rightarrow 0^+} (H - (E + i\mu))^{-1} \text{ exists in } \mathcal{L}(\mathcal{H}_1, \mathcal{H}_1^*),$$

and

$$\|u\|_{\mathcal{H}_1^*} = \|(H - (E + i0))^{-1}S\|_{\mathcal{H}_1^*} \leq c \|S\|_{\mathcal{H}_1}.$$

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$$\forall h \in]0, h_0], \quad \sup_{\substack{\operatorname{Re} z \sim E \\ \operatorname{Im} z > 0}} \|(H_h - z)^{-1}\|_{\mathcal{L}(\mathcal{H}_1, \mathcal{H}_1^*)} \leq c(h)$$

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$$\|u_h\|_{\mathcal{H}_1^*} = \|(H_h - (E + i0))^{-1} S_h\|_{\mathcal{H}_1^*} \leq c(h) \|S_h\|_{\mathcal{H}_1}.$$

We study these estimates in an **abstract setting**, and then for the **dissipative Schrödinger operator**.

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- ② We study the **semiclassical measures** for the solution u_h of the Helmholtz equation for a particular term source S_h :

$$\langle \text{Op}_{h_m}^w(q) u_{h_m}, u_{h_m} \rangle \xrightarrow{m \rightarrow \infty} \int_{\mathbb{R}^{2n}} q \, d\mu,$$

where $h_m \rightarrow 0$ and

$$\text{Op}_h^w(q)u(x) = \frac{1}{(2\pi h)^n} \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} e^{\frac{i}{h}\langle x-y, \xi \rangle} q\left(\frac{x+y}{2}, \xi\right) u(y) \, dy \, d\xi$$

(Weyl quantization of q).

Mourre's commutators method

Theorem (E.Mourre 81,...)

Let H_1 be a self-adjoint operator on the Hilbert space \mathcal{H} .

The self-adjoint operator A on \mathcal{H} is said to be *conjugate* to H_1 on the open set $J \subset \mathbb{R}$ if

- some conditions about the commutators $[H_1, iA]$ and $[[H_1, iA], iA]$ are satisfied,
- and for some $\alpha > 0$:

$$\mathbf{1}_J(H_1)[H_1, iA]\mathbf{1}_J(H_1) \geq \alpha \mathbf{1}_J(H_1).$$

In this case, for $\delta > \frac{1}{2}$ and a compact $I \subset J$ there exists $c > 0$ such that for $\operatorname{Re} z \in I$ and $\operatorname{Im} z \neq 0$

$$\left\| \langle A \rangle^{-\delta} (H_1 - z)^{-1} \langle A \rangle^{-\delta} \right\|_{\mathcal{L}(\mathcal{H})} \leq c.$$

$$\langle \lambda \rangle = (1 + |\lambda|^2)^{\frac{1}{2}}$$

Mourre's commutators method

Theorem

Let $H = H_1 - iV$ be a *dissipative* operator on the Hilbert space \mathcal{H} , where H_1 is self-adjoint and $V \geq 0$ is self-adjoint and H_1 -bounded with relative bound < 1 .

The self-adjoint operator A on \mathcal{H} is said to be *conjugate* to H on the open set $J \subset \mathbb{R}$ if

- some conditions about the commutators $[H_1, iA]$, $[V, iA]$ and $[[H_1, iA], iA]$, $[[V, iA], iA]$ are satisfied,
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Two words about the assumption

$$\mathbb{1}_J(H_1)[H_1, iA]\mathbb{1}_J(H_1) \geq \alpha \mathbb{1}_J(H_1).$$

- We do not have a **functional calculus** for the non-selfadjoint operator H .

We use functional calculus for the self-adjoint part H_1 , and the assumption that the dissipative part V is “smaller” than H_1 .

Two words about the assumption

$$\mathbb{1}_J(H_1)[H_1, iA]\mathbb{1}_J(H_1) \geq \alpha \mathbb{1}_J(H_1).$$

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We use functionnal calculus for the self-adjoint part H_1 , and the assumption that the dissipative part V is “smaller” than H_1 .

Lemma (Quadratic estimates)

Let $T = T_R - iT_I$ where T_R is self-adjoint and $T_I \geq 0$ is self-adjoint and T_R -bounded with relative bound < 1 .
If $B^*B \leq T_I$, Q is bounded and $\text{Im } z > 0$ we have

$$\|B(T - z)^{-1}Q\| \leq \|Q^*(T - z)^{-1}Q\|^{\frac{1}{2}}.$$

- We use the quadratic estimates with

$$T = H_1 - i\varepsilon\phi(H_1)[H_1, iA]\phi(H_1), \quad \text{supp } \phi \subset J,$$

$$\text{and } B = \sqrt{\varepsilon}\sqrt{\alpha}\phi(H_1).$$

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$$T = H_1 - iV - i\varepsilon\phi(H_1)[H_1 - iV, iA]\phi(H_1), \quad \text{supp } \phi \subset J,$$

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Mourre's commutators method

Theorem (J.R. 10)

Let $H = H_1 - iV$ be a *dissipative* operator on the Hilbert space \mathcal{H} , where H_1 is self-adjoint and $V \geq 0$ is self-adjoint and H_1 -bounded with relative bound < 1 .

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- and for some $\alpha > 0$, $\beta \geq 0$:

$$\mathbf{1}_J(H_1)([H_1, iA] + \beta V)\mathbf{1}_J(H_1) \geq \alpha \mathbf{1}_J(H_1).$$

In this case, for $\delta > \frac{1}{2}$ and a compact $I \subset J$ there exists $c > 0$ such that for $\operatorname{Re} z \in I$ and $\operatorname{Im} z > 0$

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More about abstract resolvent estimates

- Limiting absorption principle: for $\lambda \in J$ the limit

$$\lim_{\mu \rightarrow 0^+} \langle A \rangle^{-\delta} (H - (\lambda + i\mu))^{-1} \langle A \rangle^{-\delta}$$

exists in $\mathcal{L}(\mathcal{H})$ and defines a continuous function of λ .

- Estimate in Besov spaces.
- Estimates for the powers of the resolvent and regularity of the limit.

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Classical flow

- Let

$$H_1^h = -h^2 \Delta + V_1(x)$$

with

$$|\partial^\alpha V_1(x)| \leq c_\alpha \langle x \rangle^{-\rho - |\alpha|}, \quad \rho > 0.$$

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with

$$|\partial^\alpha V_1(x)| \leq c_\alpha \langle x \rangle^{-\rho - |\alpha|}, \quad \rho > 0.$$

- Let

$$p(x, \xi) = |\xi|^2 + V_1(x).$$

We denote by $\phi^t(x_0, \xi_0) = (\bar{x}(t, x_0, \xi_0), \bar{\xi}(t, x_0, \xi_0))$ the solution of the **hamiltonian system**

$$\begin{cases} \partial_t \bar{x}(t) = 2\bar{\xi}(t), \\ \partial_t \bar{\xi}(t) = -\nabla V_1(\bar{x}(t)), \\ \bar{x}(0) = x_0, \quad \bar{\xi}(0) = \xi_0. \end{cases}$$

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Resolvent estimates

Theorem (D.Robert-H.Tamura 87, X.P.Wang 87)

Let $\delta > \frac{1}{2}$ and $E > 0$.

Then we can find $h_0 > 0$, a neighborhood I of E and $c \geq 0$ such that for $h \in]0, h_0]$ and $\operatorname{Re} z \in I$ and $\operatorname{Im} z \neq 0$ we have

$$\left\| \langle x \rangle^{-\delta} (H_1^h - z)^{-1} \langle x \rangle^{-\delta} \right\|_{\mathcal{L}(L^2(\mathbb{R}^n))} \leq \frac{c}{h}$$

if and only if E is non-trapping:

$$p(x, \xi) = E \implies |\bar{x}(t, x, \xi)| \xrightarrow[t \rightarrow \pm\infty]{} +\infty.$$

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Resolvent estimates

Theorem (J.R. 10)

Let $\delta > \frac{1}{2}$ and $E > 0$. Suppose that $V_2 \geq 0$ is of long range.
Then we can find $h_0 > 0$, a neighborhood I of E and $c \geq 0$ such
that for $h \in]0, h_0]$ and $\operatorname{Re} z \in I$ and $\operatorname{Im} z > 0$ we have

$$\left\| \langle x \rangle^{-\delta} (H_h - z)^{-1} \langle x \rangle^{-\delta} \right\|_{\mathcal{L}(L^2(\mathbb{R}^n))} \leq \frac{c}{h}$$

if and only if for $(x, \xi) \in p^{-1}(\{E\})$

$$\sup_{t \in \mathbb{R}} |\bar{x}(t, x, \xi)| < \infty \quad \implies \quad \exists T \in \mathbb{R}, \quad V_2(\bar{x}(T, x, \xi)) > 0.$$

Resolvent estimates

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We have the limiting absorption principle and the limit of the resolvent

$$(H_h - (E + i0))^{-1} : L^{2,\delta}(\mathbb{R}^n) \rightarrow L^{2,-\delta}(\mathbb{R}^n)$$

gives the unique outgoing solution for the equation

$$(H_h - E)u = S.$$

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C. Gérard and A. Martinez (88) constructed a **conjugate operator** to H_1^h , using pseudo-differential calculus.

- We look for a conjugate operator of the form

$$A_h = \text{Op}_h^w(x \cdot \xi + r(x, \xi)), \quad r \in C_0^\infty(\mathbb{R}^{2n})$$

(if $V_1 = 0$ we can choose $r = 0$).

- In order to have

$$\mathbf{1}_J(H_1^h)[H_1^h, iA_h]\mathbf{1}_J(H_1^h) \geq c_0 h \mathbf{1}_J(H_1^h), \quad c_0 > 0,$$

after quantization, we construct r such that

$$\{p, x \cdot \xi + r(x, \xi)\} \geq c_0 \quad \text{on } p^{-1}(J).$$

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Semiclassical measure for the solution of the Helmholtz equation

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Semiclassical measure for the solution of the Helmholtz equation

Let

$$u_h = (H_h - (E + i0))^{-1} S_h$$

where

$$H_h = -h^2 \Delta + V_1(x) - ihV_2(x)$$

and S_h is an explicit source term which concentrates on a bounded submanifold of \mathbb{R}^n :

- Γ bounded submanifold of \mathbb{R}^n of dimension $d \in \llbracket 0, n-1 \rrbracket$, σ_Γ Lebesgue measure on Γ ,
- $A \in C_0^\infty(\Gamma)$,
- $S \in \mathcal{S}(\mathbb{R}^n)$,

$$S_h(x) = h^{\frac{1-n-d}{2}} \int_\Gamma A(z) S\left(\frac{x-z}{h}\right) d\sigma_\Gamma(z).$$

Semiclassical measure for the solution of the Helmholtz equation

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$$S_h(x) = h^{\frac{1-n-d}{2}} \int_\Gamma A(z) S\left(\frac{x-z}{h}\right) d\sigma_\Gamma(z).$$

We have:

$$\forall \delta > \frac{1}{2}, \quad \|S_h\|_{L^{2,\delta}(\mathbb{R}^n)} = O(\sqrt{h})$$

Semiclassical measure for the solution of the Helmholtz equation

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We have:

$$\forall \delta > \frac{1}{2}, \quad \|S_h\|_{L^{2,\delta}(\mathbb{R}^n)} = O(\sqrt{h}) \quad \text{and} \quad \|u_h\|_{L^{2,-\delta}(\mathbb{R}^n)} = O\left(\frac{1}{\sqrt{h}}\right).$$

Known results for a constant absorption index

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- J.D.Benamou-F.Castella-T.Katsaounis-B.Perthame-02:
 $\Gamma = \{0\}$, semiclassical measure as the limit of the Wigner
transform
(see also F.Castella (05)).
- F.Castella-B.Perthame-O.Runborg-02: Γ affine subspace of
 \mathbb{R}^n , $V_1 = 0$.
- X.P.Wang-P.Zhang-06: $V_1 \neq 0$.
- E.Fouassier-06: two source points.
- E.Fouassier-07: V_1 discontinuous along a hyperplane.
- J.-F.Bony-09: $\Gamma = \{0\}$, microlocal point of view.

The Assumptions

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- V_1 of long range.
- V_2 of short range:

$$|\partial^\alpha V_2(x)| \leq c_\alpha \langle x \rangle^{-1-\rho-|\alpha|}, \quad \rho > 0.$$

- E satisfies the damping assumption on trapped trajectories:

$$\sup_{t \in \mathbb{R}} |\bar{x}(t, x, \xi)| < \infty \implies \exists T \in \mathbb{R}, V_2(\bar{x}(T, x, \xi)) > 0.$$

- $\forall z \in \Gamma, V_1(z) < E.$
- if $N_E \Gamma = \left\{ (z, \xi) \in N\Gamma : |\xi|^2 + V_1(z) = E \right\}$ then

$$\sigma_{N_E \Gamma}(\{w \in N_E \Gamma : \exists t > 0, \phi^t(w) \in N_E \Gamma\}) = 0.$$

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Theorem (J.R. 10)

- *There exists a non-negative Radon measure μ on \mathbb{R}^{2n} such that*

$$\forall q \in C_0^\infty(\mathbb{R}^{2n}), \quad \langle \text{Op}_h^w(q)u_h, u_h \rangle \xrightarrow{h \rightarrow 0} \int_{\mathbb{R}^{2n}} q \, d\mu.$$

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- *μ is characterized by the following three properties:*
 - $\text{supp } \mu \subset p^{-1}(\{E\})$.*
 - $\mu = 0$ on the incoming region $\{|x| \gg 1, x \cdot \xi \leq -\frac{1}{2}|x||\xi|\}$.*
 - μ satisfies the Liouville equation*

$$\{p, \mu\} + 2V_2\mu = \underbrace{\pi(2\pi)^{d-n} |A(z)|^2 |\xi|^{-1} |\hat{S}(\xi)|^2}_{\kappa(z, \xi)} \sigma_{N_E \Gamma}.$$

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- *These three properties imply that for all $q \in C_0^\infty(\mathbb{R}^{2n})$ the integral of q is given by*

$$\int_0^\infty \int_{N_E \Gamma} \kappa(z, \xi) q(\phi^t(z, \xi)) e^{-2 \int_0^t V_2(\phi^s(z, \xi)) \, ds} \, d\sigma_{N_E \Gamma}(z, \xi) \, dt$$

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The three new difficulties:

- Non-selfadjointness of H_h .
- Geometry of Γ (and $N_E\Gamma$).
- Trapped trajectories.

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Let $w \in \mathbb{R}^{2n}$ and $q \in C_0^\infty(\mathbb{R}^{2n})$ supported close to w .

$$\begin{aligned} \text{Op}_h^w(q)u_h &= \frac{i}{h} \int_0^{T_0} \text{Op}_h^w(q) e^{-\frac{it}{h}(H_h - E)} S_h dt \\ &+ \text{Op}_h^w(q) e^{-\frac{iT_0}{h}(H_h - E)} u_h. \end{aligned}$$

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Let

$$U_1^h(t) = e^{-\frac{it}{h}H_1^h}, \quad U_h(t) = e^{-\frac{it}{h}H_h}$$

Proposition

Let $t \geq 0$ and $a \in C_b^\infty(\mathbb{R}^{2n})$. We have

$$U_1^h(t)^* \text{Op}_h^w(a) U_1^h(t) = \text{Op}_h^w(a \circ \phi^t) + O(h).$$

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$$U_h(t)^* \text{Op}_h^w(a) U_h(t) = \text{Op}_h^w \left((a \circ \phi^t) e^{-2 \int_0^t V_2 \circ \phi^s ds} \right) + O_{h \rightarrow 0}(h).$$

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Partial semiclassical measures

To avoid large times, we first study

$$u_h^T = \frac{i}{h} \int_0^T e^{\frac{it}{h}(H_h - E)} S_h dt$$

for any fixed $T \geq 0$.

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for any fixed $T \geq 0$. This gives a measure μ_T such that

$$\langle \text{Op}_h^w(q) u_h^T, u_h^T \rangle \xrightarrow{h \rightarrow 0} \int_{\mathbb{R}^{2n}} q d\mu_T.$$

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$$u_h^T = \frac{i}{h} \int_0^T e^{\frac{it}{h}(H_h - E)} S_h dt$$

for any fixed $T \geq 0$. This gives a measure μ_T such that

$$\langle \text{Op}_h^w(q) u_h^T, u_h^T \rangle \xrightarrow{h \rightarrow 0} \int_{\mathbb{R}^{2n}} q d\mu_T.$$

- $\forall \varepsilon > 0, \exists T_0 > 0, \forall T \geq T_0,$

$$\limsup_{h \rightarrow 0} |\langle \text{Op}_h^w(q) u_h, u_h \rangle - \langle \text{Op}_h^w(q) u_h^T, u_h^T \rangle| \leq \varepsilon,$$

- and

$$\int_{\mathbb{R}^{2n}} q d\mu_T \xrightarrow{T \rightarrow +\infty} \int_{\mathbb{R}^{2n}} q d\mu,$$

for some non-negative Radon measure μ on \mathbb{R}^{2n} .