## Non-Hermitian operators in QM

## \& $\mathcal{P T}$-symmetry David KREJČIǨík <br> http://gemma.ujf.cas.cz/~david/



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## Hors de ligne

(Outline)

1. QM with non-Hermitian operators (just some conceptual remarks)
2. PJ-symmetry
(what is known and my point of view)
3. physical $\mathcal{P T}$-symmetric model in QM (non-self-adjoint Robin boundary conditions)
4. Conclusions


## ¿ QM with non-Hermitian operators?



Imaginary Numbers by Yves Tanguy, 1954
(Museo Thyssen-Bornemisza, Madrid)

## Insignificant non-Hermiticity

Example 1. evolution operator $U(t)=\exp (-i t H): \quad\left\{\begin{array}{l}i \dot{U}(t)=H U(t) \\ U(0)=I\end{array}\right.$


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Theorem (spectral theorem).
Let $H=H^{*}$. Then

$$
f(H)=\int_{\sigma(H)} f(\lambda) d E_{H}(\lambda)
$$

for any complex-valued continuous function $f$.


## Technical non-Hermiticity

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Example 1. complex scaling $H_{\theta}:=S_{\theta}(-\Delta+V) S_{\theta}^{-1}, \quad\left(S_{\theta} \psi\right)(x):=e^{\theta / 2} \psi\left(e^{\theta} x\right)$


[Aguilar/Balslev, Combes 1971], [Simon 1972], [Van Winter 1974],

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Example 2. adiabatic transition probability for $H(t):=\vec{\gamma}(t / \tau) \cdot \vec{\sigma}, \quad \tau \rightarrow \infty$ [Berry 1990], [Joye, Kunz, Pfister 1991], [Jakšić, Segert 1993], . .

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Example 3. cloaking effects in metamaterials $H_{\eta}:=-\nabla \cdot a_{\eta} \nabla, a_{\eta}(x):= \begin{cases}+1, & x \in \Omega_{+} \\ -1+i \eta, & x \in \Omega_{-}\end{cases}$

[Pendry 2004], [Milton, Nicorovici 2006], [Bouchitté, Schweizer 2009], ...

# Approximate non-Hermiticity 

## open systems

## Example 1. radioactive decay



Example 2. dissipative Schrödinger operators in semiconductor physics Baro, Behrndt, Kaiser, Neidhardt, Rehberg, ...

Example 3. repeated interaction quantum systems
Bruneau, Joye, Merkli, Pillet,

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Unitary groups on a Hilbert space are generated by self-adjoint operators.

## ¿ yes ?

by changing the Hilbert space

## Non-Hermitian Hamiltonians with real spectra

$$
-\Delta+V \quad \text { in } \quad L^{2}(\mathbb{R})
$$

$$
V(x)=x^{2}+i x^{3}
$$

[Caliceti, Graffi, Maioli 1980]

$V(x)=\left\{\begin{aligned} i \operatorname{sgn}(x) & \text { if } \quad x \in(-L, L) \\ \infty & \text { elsewhere }\end{aligned}\right.$

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¿ What is behind the reality of the spectrum?


## $\mathcal{P J}$-symmetry

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\begin{gathered}
{[H, \mathcal{P T}]=0} \\
(\mathcal{P} \psi)(x):=\psi(-x) \\
(\mathcal{T} \psi)(x):=\overline{\psi(x)}
\end{gathered}
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We have in mind $H=-\Delta+V$ on $L^{2}\left(\mathbb{R}^{d}\right)$ with $\overline{V(-x)}=V(x)$.


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Here we assume that $H$ has purely discrete spectrum.


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## perturbation-theory insight

Let $H_{0}:=-\Delta+V_{0}$ be self-adjoint, with purely discrete and simple spectrum. Let $V$ be bounded and $\mathcal{P T}$-symmetric. Define $H_{\varepsilon}:=H_{0}+\varepsilon V$.
$\Longrightarrow \sigma\left(H_{\varepsilon}\right)$ is dicrete and simple $\stackrel{*}{\Longrightarrow} \sigma\left(H_{\varepsilon}\right) \cap J \subset \mathbb{R} \quad$ for every bounded $J$ and small $\varepsilon$

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## Mathematical frameworks

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\begin{array}{l|l}
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\end{array} \text { in a more general setting than: } \quad \begin{aligned}
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1. antilinear symmetry $[H, \mathcal{S}]=0$ with $\mathcal{S}$ antiunitary (bijective and $\langle\mathcal{S} \phi, \mathcal{S} \psi\rangle=\langle\psi, \phi\rangle$ ) e.g. $\mathcal{S}:=\mathcal{P T}$

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Remark. In general (in $\infty$-dimensional spaces), all the classes of operators are unrelated.

## ¿ Physical relevance ?

suggestions:

- nuclear physics [Scholtz, Geyer, Hahne 1992]
- optics [Klaiman, Günther, Moiseyev 2008], [Schomerus 2010], [West, Kottos, Prosen 2010]
- solid state physics [Bendix, Fleischmann, Kottos, Shapiro 2009]
- superconductivity [Rubinstein, Sternberg, Ma 2007]
- electromagnetism [Ruschhaupt, Delgado, Muga 2005], [Mostafazadeh 2009] experiments:
- optics [Guo et al. 2009], [Longhi 2009], [Rüter et al. 2010]


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## i but !

"So far, there have been no experiments that prove clearly and definitively that quantum systems defined by non-Hermitian PJ-symmetric Hamiltonians do exist in nature."
[Bender 2007]

## The simplest $\mathcal{P T}$-symmetric model

[D.K., Bíla, Znojil 2006]

$$
\begin{array}{ll}
\mathcal{H}:=L^{2}(0, \pi), \quad H_{\alpha} \psi:=-\psi^{\prime \prime}, ~ D\left(H_{\alpha}\right):=\left\{\psi \in W^{2,2}(0, \pi)\right. & \left.\begin{array}{l}
\psi^{\prime}(0)+i \alpha \psi(0)=0 \\
\psi^{\prime}(\pi)+i \alpha \psi(\pi)=0
\end{array}\right\} \\
\frac{-\Delta}{} \begin{array}{ll}
\frac{d \psi}{d n}-i \alpha \psi=0 & \frac{d \psi}{d n}+i \alpha \psi=0
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Theorem 1. $H_{\alpha}$ is an $m$-sectorial operator with compact resolvent satisfying

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H_{\alpha}^{*}=H_{-\alpha}=\mathcal{T} H_{\alpha} \mathcal{T} \quad(\mathcal{T} \text {-self-adjointness })
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Corollary. The spectrum of $H_{\alpha}$ is $\left\{\begin{array}{l}\text { always real, } \\ \text { simple if } \alpha \notin \mathbb{Z} \backslash\{0\} .\end{array}\right.$

## The metric operator

[D.K. 2007], [D.K., Siegl, Železný 2010]

Theorem 3. Let $\alpha \in(-1,1)$.
Then $H_{\alpha}$ is similar to a self-adjoint operator $\tilde{H}_{\alpha}:=\Theta_{\alpha}^{1 / 2} H_{\alpha} \Theta_{\alpha}^{-1 / 2}$ with

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\Theta_{\alpha}:=I+K_{\alpha}
$$

$$
K_{\alpha}\left(x, x^{\prime}\right):=\alpha e^{i \alpha\left(x-x^{\prime}\right)}\left[\tan (\pi \alpha / 2)+i \cos (\pi \alpha / 2) \operatorname{sgn}\left(x-x^{\prime}\right)\right]
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$$

Remark. As $\alpha \rightarrow 0, \quad \tilde{H}_{\alpha}=T_{\alpha}+\mathcal{O}\left(\alpha^{3}\right)$ with

$$
\begin{aligned}
\left(T_{\alpha} \psi\right)(x) & :=-\psi^{\prime \prime}(x)-\alpha^{2} \psi(x)+\frac{1}{4} \alpha^{2}[\psi(0)+\psi(\pi)] \\
D\left(T_{\alpha}\right) & :=\left\{\psi \in W^{2,2}(0, \pi) \left\lvert\, \psi^{\prime}(0)=-\psi^{\prime}(\pi)=\frac{1}{4} \alpha^{2} \int_{0}^{\pi} \psi(x) d x\right.\right\}
\end{aligned}
$$

## The physical realisation

[Hernandez-Coronado, D.K., Siegl 2010]
scattering by a compactly supported even potential $V: \quad-\psi^{\prime \prime}+V \psi=k^{2} \psi \quad k>0$


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$\underset{\text { perfect transmission }}{(\text { i.e. } R=0)} \Longrightarrow \Longrightarrow\left\{\begin{array}{c}-\psi^{\prime \prime}+V \psi=k^{2} \psi \quad \text { in } \quad(0, \pi) \\ \psi^{\prime}-i k \psi=0 \quad \text { at } 0, \pi\end{array}\right.$

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solutions given by a non-self-adjoint $\mathcal{P J}$-symmetric spectral problem:

$$
\left\{\begin{array}{rlrl}
-\psi^{\prime \prime}+V \psi & =\mu(\alpha) \psi & \text { in } \quad(0, \pi) \\
\psi^{\prime}+i \alpha \psi & =0 & \text { at } & 0, \pi \\
\mu(\alpha) & =\alpha^{2} & &
\end{array}\right.
$$

## Square-well potential


$\mu_{n}(\alpha)=\left\{\begin{array}{lll}\alpha^{2}-V_{0} & \text { if } & n=0 \\ n^{2}-V_{0} & \text { if } & n \geq 1\end{array}\right.$

perfect transmission energies: $\left\{n^{2}-V_{0}\right\}_{n=1}^{\infty}$
¿ Significance of complex spectra?
¿ Significance of complex spectra?
¡ loss of perfect transmission energies !


dispersion relations $\mu(\alpha)=\alpha^{2}$

perfect-transmission energies

# The inverse problem 

scattering data $\longrightarrow$ spectrum

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initial PTE problem

$$
\left\{\begin{array}{rlrl}
-\psi^{\prime \prime}+V \psi & =\mu(\alpha) \psi & \text { in } \quad(0, \pi) \\
\psi^{\prime}+i \alpha \psi & =0 & \text { at } \quad 0, \pi \\
\mu(\alpha) & =\alpha^{2} & &
\end{array}\right.
$$

shifted scatterer
modified initial problem

$$
\left\{\begin{array} { r l r l } 
{ - \psi ^ { \prime \prime } + ( V + V _ { 0 } ) \psi } & { = \mu _ { 0 } ( \alpha ) \psi } & { \text { in } \quad ( 0 , \pi ) } \\
{ \psi ^ { \prime } + i \alpha \psi } & { = 0 } & { \text { at } \quad 0 , \pi } \\
{ \mu ( \alpha ) } & { = \alpha ^ { 2 } } & { } & { }
\end{array} \Longleftrightarrow \left\{\begin{array}{rlrl}
-\psi^{\prime \prime}+V \psi=\mu(\alpha) \psi & \text { in } \quad(0, \pi) \\
\psi^{\prime}+i \alpha \psi & =0 & & \text { at } \quad 0, \pi \\
\mu(\alpha)=\alpha^{2}-V_{0} & &
\end{array}\right.\right.
$$

Consequently:

$$
\begin{aligned}
& \quad V_{0} \mapsto \operatorname{PTEs}\left(V_{0}\right) \Longrightarrow \alpha \mapsto \mu(\alpha) \\
& \text { (provided that } V_{0} \mapsto \operatorname{PTEs}\left(V_{0}\right) \text { are invertible) }
\end{aligned}
$$

## Conclusions

Ad $\mathcal{P T}$-symmetry:
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Generalisations:
$\rightarrow$ higher-dimensional models with both the point and continuous spectra
$\rightarrow$ curvature-induced effects
¡ many open problems! $\quad(\Leftarrow$ spectral theory of non-self-adjoint operators is "in its infancy")

ESF exploratory workshop on

# Mathematical aspects of the physics with non-self-adjoint operators 

30 August - 3 September 2010<br>Prague, Czech Republic<br>http://www.ujf.cas.cz/ESFxNSA/



Imaginary numbers 1954 by Y. Tanguy

Studying non-self-adjoint operators is like being a vet rather than a doctor: one has to acquire a much wider range of knowledge, and to accept that one cannot expect to have as high a rate of success when confronted with particular cases.
E. B. Davies 2007

## My $\mathcal{P J}$-symmetric life

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http://gemma.ujf.cas.cz/~}\mathrm{ david/
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- D.K., H. Bíla and M. Znojil: "Closed formula for the metric in the Hilbert space of a $\mathcal{P J}$-symmetric model"; J. Phys. A 39 (2006), pp. 10143-10153.
- D.K.: "Calculation of the metric in the Hilbert space of a $\mathcal{P J}$-symmetric model via the spectral theorem"; J. Phys. A: Math. Theor. 41 (2008) 244012.
- D. Borisov and D.K.: "PJ-symmetric waveguides"; Integral Equations Operator Theory 62 (2008), no. 4, 489-515.
- D.K. and M. Tater: "Non-Hermitian spectral effects in a PJ-symmetric waveguide"; J. Phys. A: Math. Theor. 41 (2008) 244013.
- D.K. and P. SiegI: "PJ-symmetric models in curved manifolds"; J. Phys. A: Math. Theor. 43 (2010) 485204.
- H. Hernandez-Coronado, D.K. and P. Siegl: "Perfect transmission scattering as a PJ-symmetric spectral problem"; submitted (2010).
- D.K., P. Siegl and J. Železný: "Non-Hermitian PJ-symmetric Sturm-Liouville operators and to them similar Hamiltonians" ; under preparation.

